

平面叶栅不可压粘性流场数值解

南京航空学院 王立成 张惠民 杨崇颖

摘 要

本文旨在发展一种二维叶栅不可压粘性流场计算的工程数值法。采用了主变量控制方程。在贴体非正交曲线坐标中求解运动方程。应用了一种新的叠加网格。用逆风因子法来确保数值解的稳定。应用压力修正代数方程来修正压力场,以满足连续方程。应用等效粘性系数紊流模型。最后对四种不同的叶栅在攻角 $-13^{\circ}\sim+11^{\circ}$ 的范围内进行了数值试验,并将落后角以及分离点的位置与平面叶栅吹风数据作了比较,相互吻合是满意的。

一 前 言

叶栅落后角的计算,对于叶背面具有分离区的流场,似乎用N-S方程求解更好一些。为了便于将计算方法推广至三维流场,本文中选用的是主变量控制方程。

首先要把叶栅通道变换到一个矩形的计算平面上去,控制方程的形式也随之要变得较为复杂。但是,引入一个有效速度以后,控制方程仍然可以保持相当简单的形式。然后,在矩形计算域内用有限差分法求解变换以后的控制方程。采用了作者早先提出的逆风因子法来确保数值解的稳定性。有关逆风因子法的细节可参看参考文献[1、2、3]。

求解不可压粘性流场的另一个困难是满足连续方程的问题。困难在于连续方程和压力 p 之间没有直接的联系。本文中建立了一种非常简单的压力修正量和连续方程之间的代数关系。采用这种方法能成倍地加快计算收敛速度。代数压力修正法的细节可参考文献[1、2]。

本文中提出了一种不同于常用的叠加网格。在这种网格中我们把速度分量 u , v 同时存储于同一个网格节点上,而把压力存储于网格单元的中心。这样,关于 u , v 的两个运动方程离散后的对应系数完全一样,从而把计算工作量差不多减少到了对一维问题的工作量。这优越性对三维问题将显得更为突出。此外,采用上述网格,使得边界条件极易施加。

综上所述,本数值法是简单、快速收敛的数值法。一个 39×19 网格点的压气机叶栅所需的CPU时间在IBM-4341计算机上约为15分钟,场内每点上的 u , v 相对误差收敛至 10^{-4} 。

用本方法和程序对四种不同的叶栅进行了数值试验。攻角范围为 $-13^{\circ}\sim+11^{\circ}$ 。将计算得到的落后角,叶背上分离点的位置等与叶栅吹风试验的结果作了比较。相互吻合是满意的。

二 计 算 方 法

1. 非正交曲线坐标系中的控制方程

不可压流非正交贴体曲线坐标系中的无因次控制方程为^[3]

$$u_{eff} \frac{\partial u}{\partial \xi} + v_{eff} \frac{\partial u}{\partial \eta} = \left[\left(\alpha \frac{\partial^2 u}{\partial \xi^2} + \gamma \frac{\partial^2 u}{\partial \eta^2} \right) / J^2 R \right] - S_x \quad (1)$$

$$u_{eff} \frac{\partial v}{\partial \xi} + v_{eff} \frac{\partial v}{\partial \eta} = \left[\left(\alpha \frac{\partial^2 v}{\partial \xi^2} + \gamma \frac{\partial^2 v}{\partial \eta^2} \right) / J^2 R \right] - S_y \quad (2)$$

$$\text{其中有效速度及源项为: } u_{eff} = (y_\eta u - x_\eta v) / J - \tau / J^2 R \quad (3)$$

$$v_{eff} = (x_\xi v - y_\xi u) / J - \sigma / J^2 R \quad (4)$$

$$S_x = \left[\left(y_\eta \frac{\partial p}{\partial \xi} - y_\xi \frac{\partial p}{\partial \eta} \right) / J + \frac{2\beta}{J^2 R} \frac{\partial^2 u}{\partial \xi \partial \eta} \right] \quad (5)$$

$$S_y = \left[\left(x_\xi \frac{\partial p}{\partial \eta} - x_\eta \frac{\partial p}{\partial \xi} \right) / J + \frac{2\beta}{J^2 R} \frac{\partial^2 v}{\partial \xi \partial \eta} \right] \quad (6)$$

$$\text{而 } J = x_\xi y_\eta - x_\eta y_\xi, \quad \alpha = x_\eta^2 + y_\eta^2, \quad \beta = x_\xi x_\eta + y_\xi y_\eta, \quad \gamma = x_\xi^2 + y_\xi^2, \quad \sigma = (y_\xi DX - x_\xi DY) / J, \\ \tau = (x_\eta DY - y_\eta DX) / J, DX = \alpha x_{\xi\xi} - 2\beta x_{\xi\eta} + \gamma x_{\eta\eta}, DY = \alpha x_{\xi\xi} - 2\beta y_{\xi\eta} + \gamma y_{\eta\eta} \quad (7)$$

上面各式中 y_ξ, y_η 等分别代表 $\partial y / \partial \xi, \partial y / \partial \eta$ 等。

由此可见, 方程(1), (2)乃是以 u_{eff} 和 v_{eff} 为对流速度的对流扩散方程, 应用逆风因子法可将其方便地离散成:

$$a_{i,j} u_{i,j} = a_{i-1,j} u_{i-1,j} + a_{i+1,j} u_{i+1,j} + a_{i,j-1} u_{i,j-1} \\ + a_{i,j+1} u_{i,j+1} - S_x \quad (8)$$

$$a_{i,j} v_{i,j} = a_{i-1,j} v_{i-1,j} + a_{i+1,j} v_{i+1,j} + a_{i,j-1} v_{i,j-1} \\ + a_{i,j+1} v_{i,j+1} - S_y \quad (9)$$

式中各系数 a 的表达式可参见文献(3)。

值得注意的是在 ξ 和 η 方向上的(8)、(9)两式中对应的 a 系数完全一样。由于数值求解中主要的计算工作量是求取系数 a , 因此, 求解(8), (9)两式实际上几乎变成求解一个运动方程。

2. 压力修正代数法

通常, 压力修正量 $\hat{p}$ 是要通过解一个关于 $\hat{p}$ 的泊松方程才能求得的[4]。本文中仅简单地用一个代数方程求解 $\hat{p}$ 。详细的公式推导可参看文献(1,2)。这里只列出要点:

(1) 将速度存放于各网格节点上, 而将压力存放于如图 1 所示的位置;

$$(2) \quad \frac{\partial p}{\partial \xi} = p_A - p_B, \quad \frac{\partial p}{\partial \eta} = p_C - p_D \quad (10)$$

$$(3) \quad p_A^{n+1} = p_A^n - \varepsilon (\nabla \cdot \vec{V}_{IE1} + \nabla \cdot \vec{V}_{IE2}) / 2 \\ p_B^{n+1} = p_B^n - \varepsilon (\nabla \cdot \vec{V}_{IE1} + \nabla \cdot \vec{V}_{IE3}) / 2 \\ p_C^{n+1} = p_C^n - \varepsilon (\nabla \cdot \vec{V}_{IE3} + \nabla \cdot \vec{V}_{IE4}) / 2 \\ p_D^{n+1} = p_D^n - \varepsilon (\nabla \cdot \vec{V}_{IE1} + \nabla \cdot \vec{V}_{IE2}) / 2 \quad (11)$$

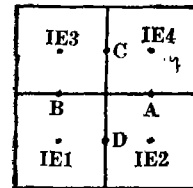


图 1 叠加网格

式中 $IE1, IE2, IE3$ 和 $IE4$ 代表如图 1 中所示的四个小单元中心上的值。 ε 是一个正值, 取在 $1/R$ 附近, 具体数值得由数值试验来确定。

如此处理压力修正, 既节省了内存又节省了运算时间, 而且无需关于 p 的任何边界条件。

3. 等效粘性系数紊流模型

采用 $k-\varepsilon$ 双方程紊流模型不仅使得程序大为复杂, 妨害程序成为一种简单的工程方法, 而且对类似于叶栅这样有曲率的通道, 也不一定能保证足够的流场预测精度^[5]。较为简单的混合长度理论在分离点上又失去了效力。本文中采用的紊流模型是等效粘性系数紊流模型^[4, 6, 7, 8], 文中同时还考虑了攻角影响的修正^[9]。

4. 边界条件与计算步骤

按本方法所进行的平面叶栅计算, 边界条件十分简单: (1) 给定叶栅进口 u, v 分布。文中假定远前方来流均匀; (2) 叶片固壁上应用无滑动条件; (3) 周期边界相应周期点上 u 和 v 分别相等; (4) 远后方出口边界上应用文献^[10]中所述的流动出口条件。求解步骤也十分简单: (1) 给定速度和压力的初场 $u^{(n)}, v^{(n)}$ 和 $p^{(n)}$; (2) 在 $p^{(n)}$ 场作用下按上述边界条件求解运动方程 (8) 和 (9) 得 $u^{(n+1)}$ 和 $v^{(n+1)}$; (3) 按 (11) 式修正压力场, 得 p^{n+1} ; (4) 转至第 (2) 步进行迭代直至收敛。

三、数值试验

对四种不同的叶栅进行了数值试验并与实验结果进行了比较。

1. 10C4/30C50叶型叶栅

该叶型及叶栅的几何参数取自文献^[11]。图 2 中画出的是进气角 $\beta_1 = 45^\circ$ 时计算得到的转角特性, 图中横坐标为攻角 i , 纵坐标为气流流过叶栅的转折角 $\Delta\beta = \beta_2 - \beta_1$ 。由图可见, 在 $-13^\circ \sim 11^\circ$ 的攻角范围的计算结果与 NACA 提供的实验数据基本上是一致的。

图 3 是同一个叶栅在 $\beta_1 = 60^\circ$ 时计算结果与 NACA 提供的实验数据和文献^[12]中无粘流计算结果的比较。可以看到当 $i > 5^\circ$ 以后, 无粘流结果与实验结果偏差较大; 而粘性流结果在攻角达 8° 时, 仍和实验结果比较地吻合。

2. J69 发动机静子末级第 III 截面串列叶栅后排叶片组成的平面叶栅 ($R = 1.3911 \times 10^5$, $\beta_y = 73^\circ$)。

图 4a 为其计算结果与南航叶栅风洞中测量得到的数据的比较。

3. 上述串列叶栅前排叶片组成的二种平面叶栅 ($R = 1.3476 \times 10^5$, $\beta_y = 55^\circ$ 和 $R = 1.3766 \times 10^5$, $\beta_y = 60^\circ$)。

图 4b 和图 4c 分别是这两种叶栅的计算和实验得到的转角特性。试验数据也是在南航叶栅

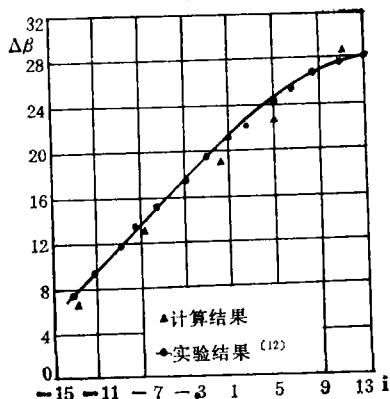


图 2 10C4/30C50 叶型叶栅转角特性 ($\beta_y = 45^\circ$)

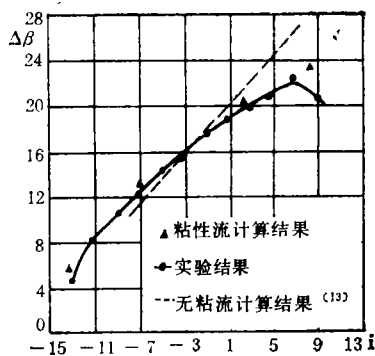


图 3 10C4/30C50 叶型叶栅转角特性 ($\beta_y = 60^\circ$)

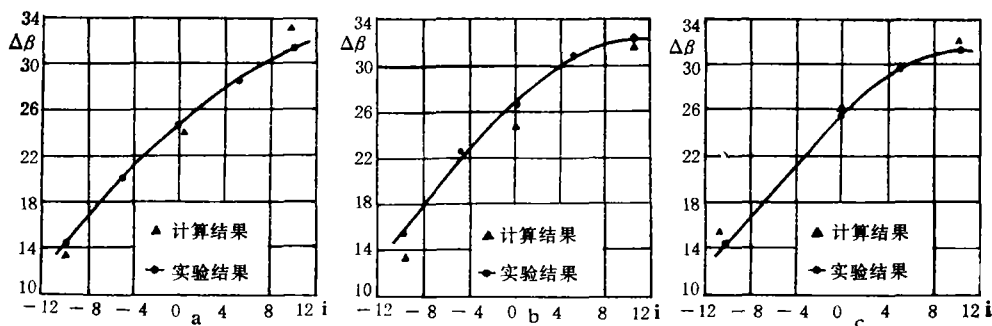


图 4 J69发动机的叶片组成的三种平面叶栅的转角特性

风洞中获得的。相互的吻合是满意的。

对叶栅中叶背分离点的位置的计算值与实验值也作了比较,在 $i = 0^\circ$ 条件下实验和计算均未发现分离。在 $i = 10^\circ$ 时,对于3.中 $\beta_y = 55^\circ$ 的叶栅,计算得到的分离点距尾缘为30mm,而叶栅涂墨试验测出的距离为28mm。对于2中的叶栅,叶背分离点距尾缘的位置计算值为11mm,实验值为15mm。

以上数值试验的结果表明,无论落后角以及分离点位置的预测都能满足工程需要。

参 考 文 献

- [1] 王立成、张惠民,“一种解对流扩散方程的有限元素逆风因子法”,航空学报, Vol.6, №5 1985。
- [2] 王立成、陈喜修,“一种简单而快速的解不可压粘性流的主变量法”,(空气动力学学报待发表)。
- [3] 王立成、张惠民,“逆风因子法在非正交曲线坐标系中解不可压N-SS方程中的应用”,航空动力学报, Vol.1, №1, July, 1986。
- [4] Busnaina, A.A. and Lilley, D.G. “A Simple Finite Difference Procedure for the Vortex Controlled Diffuser”, AIAA-82-0109, 1982。
- [5] I. Khalil and W. Tabakoff, “A Calculation Procedure for Viscous Flow in Turbomachines”, Vol. II, NASA CR 159639, and vol. III, NASA CR 159861, 1980。
- [6] Thompson, D. S., “Numerical solution of the Navier-Stokes Equations for high Reynolds Number Incompressible Turbulent Flow”, NASA-CR-163154, 1980。
- [7] Kirsnnamuthy L, “Isothermal Flowfield predictions of Confined Coflowing Turbulent Jets in an Axisymmetric Bluff-Body Near Wake”, AFWAL-TR-81-2036。
- [8] 严传俊, “模型燃烧室中流场的数值计算”, 西工大科技文献, 1982。
- [9] 杨崇颖, “平面叶栅不可压缩粘性流场的数值和实验研究”, 南航硕士研究生学位论文, 1986, 元月。
- [10] Patankar, S.V. “Numerical Heat Transfer and Fluid Flow”, Hemisphere, Washington, 1980。
- [11] Felix, A.R. and Emery, J.C., “A Comparison of Typical National Gas Turbine Establishment and NACA Axial-Flow Compressor Blade Sections in Cascade at low speed”, NACA TN 3937, 1957。
- [12] Miller, M.J., “Estimation of Deviation Angle for Axial-Flow Compressor Blade Section Using Inviscid Flow Solution”, NASA TN D-7549, 1974。

EXPERIMENTAL RESULTS OF AEROELASTIC INSTABILITY OF AN AXIAL COMPRESSOR

Feng Yucheng, Hu Zongan, Zhao Xiuhua, Wang Yuwei

(Beijing Institute of Aeronautics and Astronautics)

Abstract

This paper describes the initial experimental results of aeroelastic instability of BF-1 rotors on a transonic axial compressor facility at BIAA. Designing of BF-1 rotors depends upon the principle that the total damping is equal to zero at the point near the stall flutter boundary. The experimental results show that the design principle of BF-1 rotors is correct. Two rotors, BF-1-1 and BF-1-2, have been run from 8000 rpm to 16500 rpm. Several flow-induced vibration phenomena such as stalling flutter, wake-induced vibration, stall vibration of the blades, have been observed in experiments. In addition, a loud monotonous screaming has been heard during the stalling flutter onset. A primary theoretical analysis of this special sound is discussed. This paper also presents the experimental results of effects of the blade tip clearance and the airfoil shape on the stall flutter boundary. By comparison of the empirical methods to predict the stall flutter onset with the experimental results, it is believed that these prediction methods are acceptable.

NUMERICAL SOLUTION OF TWO-DIMENSIONAL INCOMPRESSIBLE VISCOUS FLOWS IN CASCADES

Wang Licheng, Zhang Huiming, Yang Chongying

(Nanjing Aeronautical Institute)

Abstract

This paper aims at developing an incompressible N-S equation solver for two-dimensional compressor cascade flows with massive separation for engineering applications. A staggered grid system different from that of SIMPLE method is proposed to make discretization and implementation of boundary conditions very simple and to reduce the computational efforts needed for a two- or three-dimensional problem nearly to the efforts for one-dimensional one. To ensure the numerical stability of convection-diffusion equations, a so-called upwind factor method is devised, which is robust and accurate. To deal with the determination of the pressure field an algebraic pressure correction expression is derived to simplify the numerical scheme further.

To verify the scheme four different compressor cascades are taken as examples

for numerical experiments. The incident angle varies from -13 to $+11$. The numerical results of the deviation angle and separation point are compared with the corresponding physical experimental data in good agreements.

COMPARISON OF THREE EXPERIMENTAL METHODS FOR TRANSONIC TURBINE CASCADES

Yang Xuezen

(*Gas Turbine Establishment*)

Abstract

Three typical transonic turbine cascades have been designed for this purpose and tested with the aid of three different methods (two-dimensional cascade wind tunnel test, wind tunnel test with sector of annular cascade, and hydraulic analogy). The results of these three methods are compared under the condition of geometry similarity. They are also compared with numerical solutions of the Finite Area Time Marching calculation. The strong and weak points and reliability of these methods are evaluated preliminarily.

EXPERIMENTAL INVESTIGATION OF A DOUBLE CIRCULAR ARC CASCADE WITH 73° CAMBER

Cheng Gang

(*Shenyang Aeroengine Research Institute*)

Abstract

A double circular arc (DCA) cascade with 73° camber has been tested, and according to the test results its critical Mach number, optimum incidence, deviation angle, pressure distributions on airfoil surfaces, pressure losses and incidence angle range are analysed. This cascade is also compared with a supercritical airfoil cascade. The results indicate that the DCA cascades with such large camber still can achieve good aerodynamic performances (large turning, high velocity, low loss and wide incidence angle range) and coincide well with the general principles of DCA cascades provided that the cascade parameters are chosen correctly and distribution of passage area is reasonable.

INVESTIGATION OF SUPERSONIC MULTIPLE CIRCULAR ARC (MCA) AIRFOIL CASCADE OF AXIAL-FLOW COMPRESSOR

Zha Gecheng and Yan Ruqun

(*Northwestern Polytechnical University*)

Abstract

A computational method of characteristics initial-value line in the rotational