

逆风因子法在非正交曲线坐标系中 解不可压N-S方程中的应用

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一、方法概述

所谓的逆风因子法就是分别用二阶精度公式计算出对流项和扩散项的贡献,然后按此两项贡献的比值来确定一个逆风因子,来保证对流扩散方程数值解的精度和稳定性。这种逆风因子法是作者在用有限元素法解不可压N-S方程中提出来的^[1,2]。本文将逆风因子法应用于在非正交曲线坐标系中用有限差分法解不可压N-S方程。

采用作者在文献^[2]中提出的压力场修正的代数法来满足连续方程。

非正交曲线坐标系中,不可压N-S方程的无因次形式为

$$(\vec{V} \cdot \nabla) \vec{V} = -\nabla p + \frac{1}{Re} \nabla^2 \vec{V} \quad (1)$$

$$\nabla \cdot \vec{V} = 0 \quad (2)$$

式中 $\vec{V}$ 为速度矢量,Re为雷诺数,p为静压。

在 $\xi-\eta$ 计算平面中变换以后的N-S方程可整理成如下的对流扩散方程形式:

$$u_{eff} \frac{\partial u}{\partial \xi} + v_{eff} \frac{\partial u}{\partial \eta} = \left[\left(\alpha \frac{\partial^2 u}{\partial \xi^2} + \gamma \frac{\partial^2 u}{\partial \eta^2} \right) / J^2 Re \right] - S_x \quad (3)$$

$$u_{eff} \frac{\partial v}{\partial \xi} + v_{eff} \frac{\partial v}{\partial \eta} = \left[\left(\alpha \frac{\partial^2 v}{\partial \xi^2} + \gamma \frac{\partial^2 v}{\partial \eta^2} \right) / J^2 Re \right] - S_y \quad (4)$$

其中

$$u_{eff} = \frac{1}{J} (y_\eta u - x_\eta v) - \tau / J^2 Re \quad (5)$$

$$v_{eff} = \frac{1}{J} (x_\xi v - y_\xi u) - \sigma / J^2 Re \quad (6)$$

$$S_x = \left[\left(y_\eta \frac{\partial p}{\partial \xi} - y_\xi \frac{\partial p}{\partial \eta} \right) / J + \frac{2\beta}{J^2 Re} \frac{\partial^2 u}{\partial \xi \partial \eta} \right] \quad (7)$$

$$S_y = \left[\left(x_\xi \frac{\partial p}{\partial \eta} - x_\eta \frac{\partial p}{\partial \xi} \right) / J + \frac{2\beta}{J^2 Re} \frac{\partial^2 v}{\partial \xi \partial \eta} \right] \quad (8)$$

上面各式中, u, v 为速度分量, $\alpha, \beta, \gamma, \sigma, \tau$ 和 J 等均是两个坐标系相互转换有关的系数^[3]。而 x_ξ, y_η 等代表 $\frac{\partial x}{\partial \xi}, \frac{\partial y}{\partial \eta}$ 等。

对公式 (3) 和 (4) 采用逆风因子法在计算平面内将其离散。有趣的是, 离散得到的运动方程的系数, 对 ξ 方向和对 η 方向的方程完全一样, 这样, 求解两个运动方程实际上减少到了求解一个运动方程。这是使本方法收敛快速的原因之一。

二、数值试验

为了验证本文中提出的计算方法, 对空腔和突然扩张通道中的流场作了数值经验, 并与对应的其他数值方法和物理实验结果作了比较。

1. 空腔流场的数值试验

试验雷诺数为 $Re = 100.0$ 。图 1 是按本文方法计算出来的流场图。而图 2 中是文献(4)中的数值结果。相互吻合相当好。

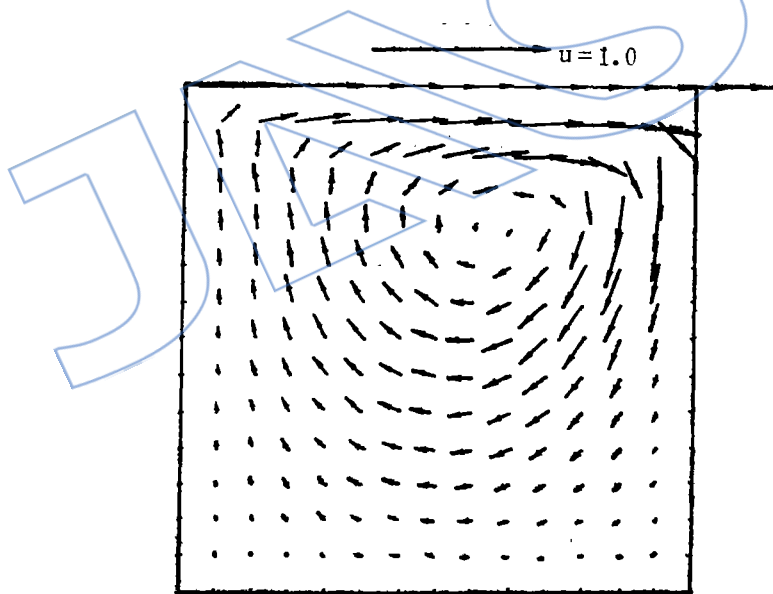


图 1 $Re = 100.0$ 时空腔内计算得到的速度矢量分布

2. 突然扩张通道中的流场计算

突然扩张通道的物理实验数据取自文献(5), 并且所取的通道几何尺寸完全相同。图 3 中所示的是计算得到的 u 速度分布曲线, 分离区用虚线与非分离区相区分。而图 4 是对应的

物理试验数据。对比图 3 和图 4, 发现无论是速度分布形状还是分离区的形状和大小相互吻合很好。从文献〔5〕中可以查到当 $Re_h = 73$ 时 (Re_h 为以 h 为参考长度的雷诺数, 本文中的计算也取 $Re_h = 73$), 分离区长度 $x/h \approx 4.4$, 而本文计算出来的分离区长度为 $x/h \approx 4.6$ 。

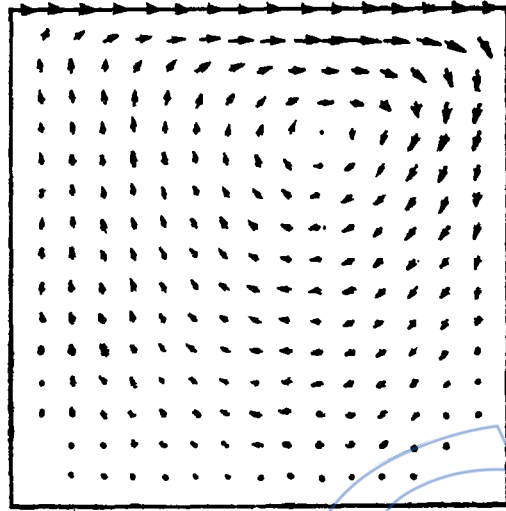


图 2 按文献〔4〕中的数值法计算得到的空腔内的速度矢量分布图

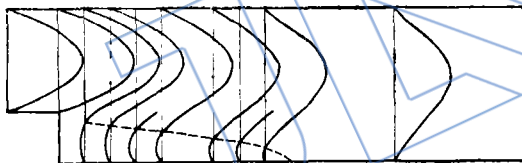


图 3 按本文方法计算得到的突然扩张通道内 u 速度分量的分布图

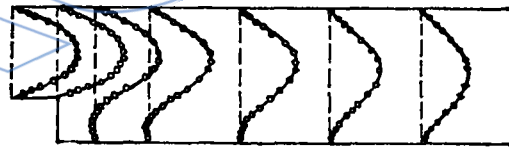


图 4 突然扩张通道内 u 分量实验值〔5〕

参 考 文 献

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the curvature of the surface and the presence of the ejection holes are in favour of the heat transfer augmentation. It is noted that the internal impingement heat transfer coefficient on the suction side of the cooled blade calculated by using the correlation of that on flat plate is more or less too conservative.

APPLICATION OF THE UPWIND FACTOR METHOD TO SOLUTION OF INCOMPRESSIBLE NAVIER-STOKES EQUATIONS IN NON-ORTHOGONAL CURVILINEAR COORDINATE SYSTEM

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Abstract

The upwind factor method has found successful application in the finite element scheme. The purpose of this paper is to extend its application to solution of incompressible Navier-Stokes equations in non-orthogonal curvilinear coordinate system by the finite difference method. An effective velocity is introduced to keep the control equations in the computational domain as simple as in the physical domain.

CORRESPONDENCE

DIMENSIONLESS ANALYSIS OF LIQUID IMPINGING PLANE FILM

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Abstract

Seven nondimensional parameters which describe the liquid phase mixing phenomenon of a pair of liquid plane impinging streams produced by two plane slot injectors are determined under specific conditions. The results provide an analytical means for solving liquid phase mixing problems by experiments